

### MAGNETIC HARDENING INDUCED BY EXCHANGE COUPLING IN MECHANICALLY MILLED ANTIFERROMAGNETIC – FERROMAGNETIC COMPOSITES

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### ABSTRACT

Mechanical milling has been used to synthesize ferromagnetic (FM, Co) - antiferromagnetic (AFM, NiO) composites. The coercivity, H<sub>C</sub>, and energy product, BH<sub>Max</sub>, of these composites can be enhanced at room temperature after appropriate heat treatments above the Néel temperature of the AFM, T<sub>N</sub>. Although the maximum H<sub>C</sub> is achieved for the (NiO)1:1(Co) weight ratio, BH<sub>Max</sub> is further enhanced for the (NiO)2:3(Co) ratio, where higher saturation magnetization is obtained due to the larger amount of FM. Exchange coupling, responsible for these effects, decreases as the temperature is increased and vanishes close to T<sub>N</sub>. The thermal stability of the coercivity enhancement remains rather insensitive to the somewhat broad distribution of blocking temperatures of this system.

### INTRODUCTION

Since mechanical alloying (MA) was developed in the late 1960s [1], this technique has been used as a processing route for the synthesis of a large variety of equilibrium and non-equilibrium phases and phase mixtures [2]. This technique has been shown to be successful in reducing crystallite sizes to the nanometric range. These small particles can exhibit rather different physical properties from those of their bulk counterpart [3].

Mechanical alloying has also been used to produce direct or indirect magnetic interactions, like in exchange spring magnets (milling soft-hard ferromagnets) [4,5] or giant magnetoresistance materials (milling magnetic-nonmagnetic phases) [6]. However, studies of exchange interaction between ferromagnetic and antiferromagnetic materials induced by ball milling are scarce. Nevertheless, recently the possibility of increasing the room temperature coercivity by mechanical alloying transition metals (Ni, Co, Fe) and their own antiferromagnetic oxides (NiO or CoO) or sulfides (FeS) has been proven [7]. Moreover, it is well known that oxidized ferromagnetic fine particles can also exhibit coercivity enhancements [8]. However, this enhancement is mainly observed far below room temperature, which makes this property not useful for applications. It is noteworthy that coercivity increases due to AFM-FM exchange coupling are also being studied in thin film systems [9,10].

In this paper we show the possibility of increasing the room temperature coercivity and

In this paper we show the possibility of increasing the room temperature coercivity and energy product by ball milling FM and AFM composed of different transition metals, e.g. Co and NiO, after adequate magnetic field heat treatments. The hardening of the FM phase is optimized by varying the AFM:FM ratio.

## EXPERIMENTAL

Different AFM:FM weight ratios (0:1, 3:7, 2:3, 1:1 and 3:2) of gas - atomized powders of NiO (99%, <100  $\mu$ m,  $T_N$  = 590 K) and Co (99.5%, < 44  $\mu$ m,  $T_C$  = 1404 K) were mixed and sealed under Ar atmosphere in an agate vial (V = 20 ml), together with 6 agate balls (10 mm diameter), with a ball to powder weight ratio of 2:1. The milling was carried out for different

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times (0.1-30 hours), at 500 rpm, using a planetary mill. To induce exchange coupling, the as milled powders were annealed at different temperatures, TANN = 300 - 850 K, for different times (0-5 hours) under vacuum, in the presence of different magnetic fields (0 - 10 kOe), and field cooled to room temperature.

The microstructure of the powders was studied using standard x-ray diffraction (XRD) (Cu-K\approx radiation). The x-ray diffraction peaks were deconvoluted and fitted using a pseudo-Voigt function. The crystallite size, \(\sigma\rightarrow\ri

# RESULTS AND DISCUSSION

Shown in fig. 1(a) is a SEM image (secondary electron) for the NiO-Co powders ball milled for 20 hours in a weight ratio of 1:1. Ball milling induces a broad range of particle sizes, from 1  $\mu m$  to 30  $\mu m$  [11]. The larger particles (an enlargement of one of these particles is shown in fig. 1) are composed of about 1  $\mu m$  thick lamellae, surrounded by small particles. As can be seen in fig. 1(b) the Co x-ray mapping indicates that the lamellae (bright in fig. 1(a)) correspond to Co while the particles (gray in fig. 1(a)) correspond to NiO. This microstructure, typical of ball milled ceramic-metal mixtures, is due to the ductile and brittle character of Co and NiO respectively. Moreover, this microstructure allows the existence of a large interface area between the FM and the AFM phases.

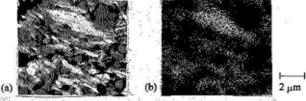


Fig. 1 (a) SEM secondary electron image and (b) Co EDX mapping for a (NiO)1:1(Co) ball milled for 20 h.

X-ray diffraction (XRD) patterns have been obtained for the unmilled Co and NiO and for the 20 h ball milled NiO-Co in a weight ratio of 1:1, before and after field annealing them in a 5 kOe magnetic field at Tann = 600 K. The XRD patterns of the unmilled Co reveal that it corresponds to a mixture of hcp and fcc Co, while the unmilled NiO powder diffraction peaks correspond to a fcc phase. The crystallite sizes, evaluated for the Co hexagonal (002) and the NiO (111) peaks, are <D>Co, blex = 42 ± 4 nm and <D>NiO > 100 nm, respectively [11].

When Co and NiO powders are ball milled together in a 1:1 weight ratio the intensity of the peaks corresponding to the fcc Co phase rapidly decreases. This allotropic transformation has been previously reported for ball milled Co powders [12]. The peak widths of the different phases after ball milling show an important increase. This is related to the crystallite size reduction, the introduction of planar defects and the microstrain increase. The crystallite sizes, after 20 hours of milling, estimated from the (002)co,Hex and the (111)NiO reflections, are: <D>co,Hex = 10 ± 2 nm and <D>nio = 14 ± 2 nm, respectively. In addition, the positions of the

NiO peaks do not change after ball-milling, indicating no important variation in its lattice parameter. This suggests that the diffusion of Co into the NiO cell is very small [11].

When the ball milled NiO-Co powders are annealed at Tann = 600 K for 0.5 hours the x-ray diffraction pattern is very similar to that of the as milled powders, where the NiO peaks are not displaced and only a minor decrease in the peak width is obtained. The crystallite size calculated from the (111)no reflection is <D>no = 18 ± 2 nm. The Co-hop phase shows a similar behavior, with <D>no = 18 ± 2 nm, calculated from the (002)co,no peak. Moreover, at this same aling temperature, a slight increase in the intensity of the (200)co,no peak can already be observed, indicating the starting point for an allotropic transformation from hep to hep-fee [12]. However, after annealing at Tann = 740 K for 0.5 hours, the crystallite sizes, calculated from (100)co,no and (111)no peaks, increase significantly to <D>co,los = 44 ± 4 nm and 4D>no = 40 ± 4 nm and the intensity of the peaks corresponding to the cubic phase increases. Thus, annealing at Tann = 740 K induces an important crystal growth and the allotropic phase transformation from hep to fee Co [11].

The room-temperature coercivity, Hc, has been measured as a function of the milling time for the different NiO:Co weight ratios for both as-milled and annealed (Tabet = 600 K) powders. In pure Co, He increases after a short milling time (1 hour) from Hc = 170 Ge (as obtained) to 290 Ge (1 hour bell milled) [13]. This increase is stributed to the allotropic phase transformation from Soc+hop Co to hop Co, where hep Co is known to have larger He than for Co [14]. Further milling decreases He (e.g. Hc = 125 Ge after milling 30 hours), which is probably linked with the disordering of the hexagonal structure, as has been reported for long-term milling [12]. A maximum in He is also observed for ball milled NiO:Co. However, for each AFM:FM ratio, the largest He is obtained for different milling times, which increase as the NiO content increases, indicating that NiO probably delays the Co structural changes.

Annealing the ball milled pure Co powders at Taba = 600 K (i.e. Tn < Taba < Tc) for 0.5 hours in H = 5 kOe results in a slight decrease in Hc. However, a significant enhancement of Hc is observed after annealing the ball milled NiO-Co powders at the same temperature, Taba = 600 K, and field cooling (H = 5 kOe) to room temperature [13]. Shown in fig. 2 is the dependence of the coercivity enhancement (defined as the difference of Hc after and before annealing at Taba, Alic) on the annealing temperature for (NiO)1:1(Co) ball milled 20 h. As can be seen in the figure, the room-temperature Hc increases only when the annealing temperature is larger than the NiO Néel temperature (i.e. Tr = 590 K), as expected for an exchange coupled FM-AFM system [9]. The decrease of AHc after annealing at exceedingly high temperatures (Taba) > 700 K) is due to the allotropic phase transformation from fee+hep Co to fee Co, as

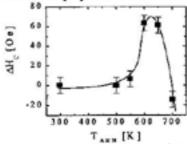


Fig. 2 Dependence of the coercivity enhancement, ΔH<sub>C</sub>, on the annealing temperature for a (NiO)1:1(Co) ball milled for 20 h.

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observed from XRD. Loop shifts in the field axis, Hg, towards the negative field axis are also present in the hysteresis loops of the ball milled NiO-Co samples for all AFM:FM ratios, except in pure Co. The loop shifts depend on the annealing temperature with a similar dependence as  $\Delta$ Hc, confirming the assumption that the exchange coupling between the AFM and the FM is responsible for both effects,  $\Delta$ Hc and Hg.

Exchange coupling produces not only an enhancement of Hc but also an enhancement of the squareness, Ms/Ms [13], and consequently an increase of the energy product, BH<sub>Max</sub>. The dependence of He, ΔHc and ΔBH<sub>Max</sub> (defined as the difference of BH<sub>Max</sub> after and before annealing at T<sub>ANN</sub> = 600 K) on the Co content (NiO:Co ratio) is shown in fig. 3. The values shown in the figure correspond to the milling times which give the maximum values of Hc for each NiO:Co ratio. As can be seen in the figure, the largest exchange coupling effects, i.e. maximum enhancements of Hc and BH<sub>Max</sub> and maximum loop shifts, are exhibited for the (NiO)1:1(Co) ratio. However, the maximum value of BH<sub>Max</sub> after annealing at T<sub>ANN</sub> = 600 K is not obtained for (NiO)1:1(Co) (BH<sub>Max</sub> = 62000 G.Oe) but for (NiO)2:3(Co) (BH<sub>Max</sub> = 83000 G.Oe). This is due to the interplay between Hc and Ms on BH<sub>Max</sub>. The addition of an AFM and the subsequent heat treatments increase Hc, however the AFM also promotes a reduction in the overal! Ms of the composite. Hence, the enhancement of BH<sub>Max</sub> is a compromise between both effects.

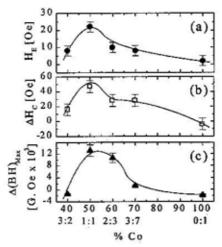


Fig. 3 Dependence of the (a) exchange bias, H<sub>E</sub>, (b) coercivity enhancement, ΔH<sub>C</sub>, and (c) enhancement of the energy product, Δ(BH)<sub>Max</sub>, on the Co percentage (NiO:Co ratio).

Shown in fig. 4 is the dependence of the loop shift on the measuring temperature. As expected from exchange coupled AFM-FM systems, He decreases as measuring temperature increases, becoming zero at  $T \approx 600$  K, i.e. at the Néel temperature of the AFM (Tn (NiO) = 590 K). A similar behavior is observed for  $\Delta$ Hc and  $\Delta$ (BH)<sub>Max</sub>. The decrease of these properties with increasing measuring temperature is the result of a reduction of the AFM/FM interface coupling,

as a consequence of either the loss of the magnetic order in the AFM and/or the reduction of the AFM magnetocrystalline anisotropy. The critical temperature at which exchange coupling effects completely venish is usually denoted as the blocking temperature, Ts [9].

To study the thermal stability of the ball milled NiO:Co systems due to the concomitant distribution of blocking temperatures [15] (brought about by the distribution of crystallite sizes and defects), the NiO:Co samples were warmed to a temperature To, 300 K < To < 700 K, and subsequently cooled to room temperature in a field H = - 5 kOe, i.e. opposite to the one used to induce the original coupling (H = 5 kOe).

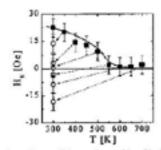


Fig. 4 Temperature dependence of the exchange bias, H<sub>0</sub> (\*), on the measuring temperature, and H<sub>0</sub> (P) after field cooling in H = -5 kOe from different temperatures.

Following this procedure, if the local blocking temperature of a crystallite (Tsi) is below To, the crystallite becomes paramagnetic, thus losing the information of the original coupling. Cooling such crystallite in H = -5 kOe induces a coupling opposite to the original one, hence shifting the loop towards positive fields. However, crystallites with Tisi > To remain unchanged. Consequently, the Hz obtained after warming to each To gives a measure of the number of crystallites with Tisi < To [15]. As can be seen in fig. 4, already at To = 350 K a reduction in Hz can be observed. Moreover, for To = 500 K Hz is reduced to zero (i.e. 50 % of the crystallites have  $T_{\rm Bi} < 500$  K). These results indicate a broad distribution of blocking temperatures in this system. However, although the thermal stability of Hz is strongly influenced by the distribution of Ta, we observed that Hz is rather intensitive to it. This is probably because although the overall Hz depends on the strength and sign of the coupling for each crystallite, the enhancement of Hc depends only on the strength of the coupling.

## CONCLUSIONS

In conclusion, we have shown that mechanically milled AFM(NiO) - FM(Co) composites display significant enhancements of the room temperature coercivity and energy product due to the exchange coupling induced after heat treating the samples at Toos > To under magnetic fields. The as-milled powders exhibit a peculiar microstructure (Co lamellac surrounded by refined NiO particles), which allows the existence of a large interface area between the FM and the AFM. The maximum magnetic hardening is shown to be a trade off between the coercivity enhancement produced by exchange coupling and the overall reduction of saturation magnetization due to the antiferromagnet. All the effects (Hc, He and BHose) decrease with increasing temperature and the NiO-Co samples become uncoupled when the

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measuring temperature is larger than the antiferromagnetic Néel temperature. Although  $H_{\rm E}$  is strongly influenced by the distribution of  $T_{\rm B}$ ,  $H_{\rm C}$  remains rather insensitive to it.

### **ACKNOWLEDGEMENTS**

This work was supported in part by DGESEIC under contract MAT98-0730 and by CICYT under contract PETR195-0311-OP. J.S. thanks the DGU for his fellowship. J.N. and X.A. acknowledge the Spanish Government for its financial support.

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